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 $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ Mesa

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Polarization Enhancement of terahertz radiation generated by intrinsic Josephson junctions in a truncated edge square $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ mesa.

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Abstract

In this study, we investigated the terahertz radiation from a truncated edge square mesa structure made from a superconducting $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$. Using a commercial software, the polarization characteristics were determined, and introduced, while accounting for the skin effect. The axial ratio was enhanced in the simulation by performing a parametric study on the design.

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Keywords: intrinsic Josephson junctions; THz emission; polarization

1. Introduction

In the last decade, great advances have been made in the field of Terahertz (THz) continuous-wave sources. THz radiation sources and methods have increased in number and improved in efficiency and power [1]. One of the most recently discovered THz sources is built from the high- T_c superconductor $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ (BSCCO) [2,3]. Utilizing the ac Josephson effect, the radiation in BSCCO based emitters is generated by an oscillating current in the intrinsic

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Josephson junctions (IJJs). The IJJs are formed as superconducting CuO_2 layers, stacked over insulating Bi-Sr-O layers along the c-axis of a mesa structure. Electromagnetic waves in general can be classified into three main classifications according to their polarization, linear, elliptical, and circular. However, only linearly polarized THz radiation was observed from BSCCO emitters so far [4]. Many new experiments and applications could be proposed if the polarization can be controlled in such emitters. Examples may include vibrational circular dichroism (VCD) spectroscopy, high speed communications and THz continuous-wave polarization imaging.

According to previous studies [5][6], the mesa structure resembles a microstrip patch antenna, therefore, Antenna Theory can be used to analyse the measured radiation properties, and — as in our study — predict it. The excitation of resonant modes plays a vital role in the emission of THz radiation from BSCCO mesa. Two spatially orthogonal resonant modes, which are equal in amplitude and differ in phase by 90 degrees, have to be excited in the mesa to observe a circularly polarized (CP) radiation. These two modes are resonant along the diagonal axis of the mesa, and can be controlled by modifying the areas of the truncated part Δs and square space S of the mesa as described by these relations [7]:

$$f_{0r} = \frac{c_0}{2nw}, \quad f_a = f_{0r}, \quad f_b = f_{0r} \left(1 + \frac{2\Delta s}{S}\right)$$

Where c_0 , n , w , are the speed of light in vacuum, the refractive index, and the width of square mesa, respectively. The main design parameters for the truncated edge square mesa that has a side length A are, a_1 the length of the truncated edge, and a_2 the length of the uncut part as shown in Figure 1, where $A = a_1 + a_2$, $S = A^2$ and $\Delta s = a_1^2$. Note that f_{0r} , f_a , f_b are the resonant frequency of the uncut patch, and the frequencies of the two new orthogonal modes of truncated patch, respectively. The condition for CP radiation occurs at the arithmetic mean value of these two resonant frequencies.

In this paper, the polarization characteristics of a truncated corner mesa design are introduced. This novel structure was used in another experiment, to generate an elliptically polarization radiation from a BSCCO THz emitter, which was the first observation of a non-linear polarization from a BSCCO THz source. The preliminary data in the experiment showed a highly elliptical polarization with an axial ratio approaching to 4 dB in the frequency range predicted by the cavity model. The details will be described elsewhere [8]. To enhance the polarization, an

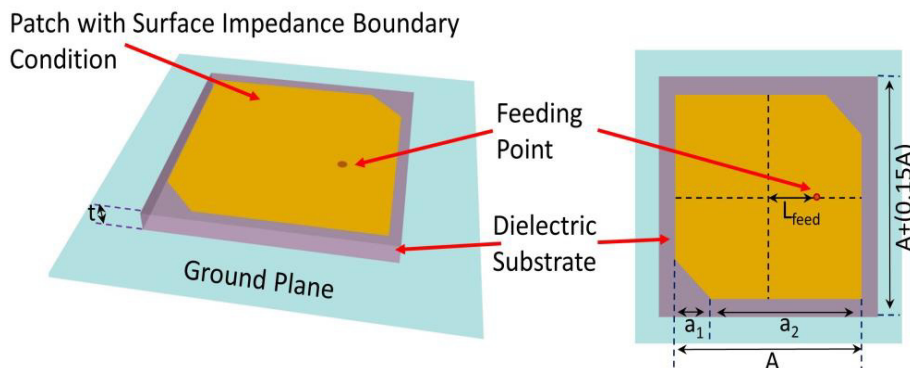


Fig. 1. A schematic view of the simulation model used in this study. The figure in the right shows a top view for the model.

electromagnetic simulation (EM) is used to determine the axial ratio (AR) of emitted waves. Then by performing a parametric study, the enhanced axial ratios is observed and compared to the measured ones.

2. Simulation Method

A commercial EM simulator named High Frequency Structural Simulator (HFSS) from ANSYS, Inc. was used to perform the simulation in this study. Using finite-element method (FEM), HFSS solves Maxwell's equation in 3D to find the required electromagnetic characteristics of a model. A similar method was used by Tsujimoto et al [9] using Sonnet EM simulator to identify the cavity modes of a nearly square mesa structure. Losses due to the frequency dependent skin effect was accounted for by applying a frequency dependant impedance boundary condition [10,11]. The complex surface impedance $Z_s = R_s + iX_s$, where R_s is the surface resistance, and X_s is the surface reactance,

was found to have a significant effect on the electromagnetic properties of the mesa. The real part of the impedance R_s corresponds to the density of the unpaired charge carriers i.e., the conduction loss. In the other hand the imaginary part X_s is associated with the density of the paired charge carriers, which relates to the penetration depth $\lambda_L(T)$, and also describes the frequency dependent phase shift of the conductor. Layered superconductors like BSCCO contain very thin layers of superconductors separated by thicker dielectrics. Due to this layered nature, the current value and nature in the c-axis and ab-plane are quite different, where in the ab-plane it is similar to the bulk homogeneous superconductors, it can be described in the local London limits [12]. To estimate the surface impedance in this case, a self-contained computer code proposed by Zimmermann et al [11,13] was used. Despite being very important to explain the EM emission from IJJs, in this paper we didn't take IJJs synchronization into account. That being said, the simulation method used in here shows a good agreement with the experiments [8], and might help to clarify some radiation characteristics and improve the polarization.

The model used in the simulation is shown in Figure 1, in which two thin layers of perfect conductors are set over and beneath a dielectric box with a dielectric constant of $\epsilon = 17.6$. The dimensions for the simulation model used in here, are correlated with those of a typically emitting square mesa structure [14].

3. Results and Discussions

Using the method mentioned earlier, the frequency dependent real part of the surface impedance was determined to be in the range $R_s = 0.159 - 0.15906 \Omega$ in a frequency range of $f = 420 - 440$ GHz. The imaginary part of the surface impedance was in the range $X_s = 1.34 - 1.4 \Omega$ within the same frequency range. The material parameters used to perform the simulation are of a typical BSCCO THz emitting mesa, in which the transition temperature $T_c = 82$ K, the emission temperature $T_b = 40$ K, the mean free path $l = 10$ nm, the London penetration depth $\lambda_L = 400$ nm, and the coherence length $\xi_0 = 0.1$ nm. In general, the radiation characteristics acquired by applying the surface impedance boundary condition as performed in this study are in agreement with the experimental study, in contrary to the simulation done without applying it, which was not correlated to experimental results. The polarization direction in simulation was found to be a left hand polarization, where in the experiments it was not determined.

To enhance the polarization in similarly shaped mesa structures, two parameters were considered in the simulation, a_1 and a_2 . At first a_2 was set to a fixed value of $68.7 \mu\text{m}$, and the a_1 's value was incremented by $1 \mu\text{m}$ in every next run, starting with $a_1 = 5 \mu\text{m}$. In figure 2 (b), the values for axial ratio (AR) for two values of a_1 swept in the frequency range of $f = 420 - 490$ GHz as determined from the simulation is shown. The red line represents the AR at $a_1 = 13 \mu\text{m}$. At this value of a_1 , the AR was elliptically polarized with a peak at $f = 440$ GHz, though its value didn't exceed 5.3 dB. An enhanced value for the AR was found when a_1 was set to $9 \mu\text{m}$, as shown by a black line in the figure. The minimum AR at this a_1 value was 0.77 dB at a frequency of 459 GHz, which is about 85.5% enhancement. At the minimum AR, the ratio between the truncated length a_1 to the uncut part length a_2 was at 13.1%. By decreasing the value of a_1 from $13 \mu\text{m}$ to $9 \mu\text{m}$, the AR peak frequency has shifted by 19 GHz to the right (increased), yet it is still within the range attainable by experimentally produced IJJs emitters. The axial ratio band width (ARBW) at 3 dB was improved from 0 to 4 GHz (460.6 GHz - 456.6 GHz). In figure 2 (a), the scattering parameter is presented using the same colours explained for the two values of a_1 . The impedance bandwidth (ZWB) at -10 dB was also improved from 0 to 2.8 GHz (465.5 GHz - 462.7 GHz).

Using a similar method, the second parameter a_2 was altered while the parameter a_1 was fixed to $13 \mu\text{m}$. The results are shown in figure 2 (c)(d), where in figure 2 (d) the axial ratios for the values $a_2 = 68 \mu\text{m}$, and $94 \mu\text{m}$ are illustrated by a red line, and a black line, respectively. This time by increasing the a_2 value, an enhancement of 93.6% is realized, where the minimum AR was 5.6 dB at the first value of a_2 , and 0.36 dB after enhancement. The frequency shift in AR peak was much larger than before, where it has shifted to the left (decreased) by 113 GHz from 444 GHz to 331 GHz. Comparing to the last parameter ratio value, a relatively similar a_1 to a_2 ratio of 13.8% was required to achieve the minimum AR. The ARBW has improved from 0 GHz in the case of $a_2 = 13 \mu\text{m}$ in the simulation to 3.62 GHz (332.92GHz - 329.3GHz). Figure 2 (c) shows the scattering parameter values in the same frequency range. A very large frequency shift and peak difference values are also observed in it, though the peaks did not pass the -10 dB range.

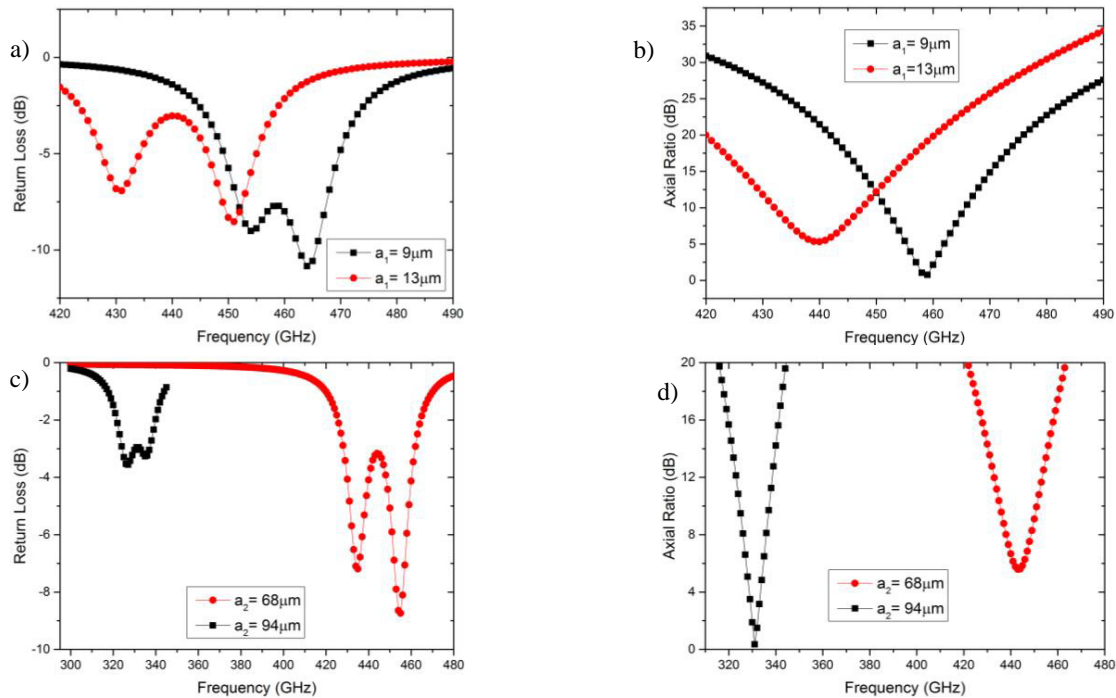


Fig. 2 (a) scattering parameter vs. frequency in the case of $a_2=68.7 \mu\text{m}$, $a_1=9, 13 \mu\text{m}$; (b) the axial ratio for the mentioned values of a_2, a_1 ; (c) scattering parameter when $a_1=13 \mu\text{m}$, $a_2=68, 94 \mu\text{m}$ (d) the axial ratio in this case.

4. Conclusion

In this study, the polarization characteristics of a truncated edge square mesa are numerically simulated using HFSS and impedance boundary condition. By using a parametric study approach, an enhancement for the polarization is realized, and discussed.

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References

- [1] M. Tonouchi, Cutting-edge terahertz technology, Nat. Photonics. 1 (2007) 97–105.
- [2] L. Ozyuzer, et al., Emission of coherent THz radiation from superconductors., Science. 318 (2007) 1291–1293.
- [3] U. Welp, K. Kadowaki, R. Kleiner, Superconducting emitters of THz radiation, Nat. Photonics. 7 (2013) 702–710.
- [4] H. Minami, I. Kakeya, et al, Characteristics of terahertz radiation emitted from the intrinsic Josephson junctions in high- T_c superconductor $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$, Appl. Phys. Lett. 95 (2009) 232511.
- [5] K. Delfanazari, et al., Tunable terahertz emission from the intrinsic Josephson junctions in acute isosceles triangular $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ mesas, Opt. Express. 21 (2013) 2171.
- [6] K. Kadowaki, et al., THz Wave Emission from Intrinsic Josephson Junctions in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$, J. Phys. Conf. Ser. 400 (2012) 022041.
- [7] R.A. Sainati, CAD of microstrip antennas for wireless applications, Artech House, (1996).
- [8] A. Elarabi, Y. Yoshioka, M. Tsujimoto, Y. Nakagawa, I. Kakeya, to be published.
- [9] M. Tsujimoto, Cavity mode Identification for coherent terahertz emission from High- T_c superconductors, arXiv: 1509.02985 (2015).
- [10] A.R. Kerr, Surface impedance of superconductors and normal conductors in EM simulators, Electron. Div. Intern. Rep. (1999) 1–16.
- [11] K.K. Berggren, D.P.L. Aude, Modeling superconductors using surface impedance techniques, (2010).
- [12] T.M. Slipchenko, et al., Surface & waveguide Josephson plasma waves in slabs of layered superconductors, Phys. Rev. B. 84 (2011) 224512.
- [13] W. Zimmermann, et al., Optical conductivity of BCS superconductors with arbitrary purity, Phys. C Supercond. 183 (1991) 99–104.
- [14] M. Tsujimoto, et al., Geometrical resonance conditions for THz radiation from the intrinsic Josephson junctions in $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$, Phys. Rev. Lett. 105 (2010) 16–19.